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Kontakt/Contact

[Digizeitschriften e.V.](#)
SUB Göttingen
Platz der Göttinger Sieben 1
37073 Göttingen

✉ info@digizeitschriften.de

HARMONIC MAPPINGS OF SURFACES

ALOIS ŠVEC, Olomouc

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We are going to study the harmonic and slightly less than harmonic mappings $f: M \rightarrow N$ in the case $\dim M = \dim N = 2$. For further details, see [1]–[9].

1. Let M, N be Riemannian manifolds, $\dim M = \dim N = 2$, $f: M \rightarrow N$ a mapping; everything be of class C^∞ . Let us suppose that M and N are oriented and $f: M \rightarrow N$ is orientation preserving. Let M be covered by a system of domains such that in each of them we are able to choose a field of orthonormal frames $\{v_1, v_2\}$, let $\{\omega^1, \omega^2\}$ be the dual bases. The Euclidean connection of M is then given by

$$(1.1) \quad \begin{aligned} \nabla m &= \omega^1 v_1 + \omega^2 v_2, \quad \nabla v_1 = \omega_1^2 v_2, \quad \nabla v_2 = -\omega_1^2 v_1; \\ d\omega^1 &= -\omega^2 \wedge \omega_1^2, \quad d\omega^2 = \omega^1 \wedge \omega_1^2, \quad d\omega_1^2 = -K\omega^1 \wedge \omega^2, \end{aligned}$$

K being the curvature of M . Analogously, the connection on N be given by

$$(1.2) \quad \begin{aligned} \nabla^* n &= \Omega^1 v_1^* + \Omega^2 v_2^*; \quad \nabla^* v_1^* = \Omega_1^2 v_2^*, \quad \nabla^* v_2^* = -\Omega_1^2 v_1^*; \\ d\Omega^1 &= -\Omega^2 \wedge \Omega_1^2, \quad d\Omega^2 = \Omega^1 \wedge \Omega_1^2, \quad d\Omega_1^2 = -K^*\Omega^1 \wedge \Omega^2. \end{aligned}$$

On M , we get the induced forms

$$(1.3) \quad \tau^1 := f^*\Omega^1, \quad \tau^2 := f^*\Omega^2, \quad \tau_1^2 := f^*\Omega_1^2$$

satisfying

$$(1.4) \quad d\tau^1 = -\tau^2 \wedge \tau_1^2, \quad d\tau^2 = \tau^1 \wedge \tau_1^2, \quad d\tau_1^2 = -K^*\tau^1 \wedge \tau^2.$$

Let us write

$$(1.5) \quad \tau^1 = a_1\omega^1 + a_2\omega^2, \quad \tau^2 = a_3\omega^1 + a_4\omega^2.$$

Then

$$(1.6) \quad \tau^1 \wedge \tau^2 = \mu\omega^1 \wedge \omega^2, \quad \text{where } \mu = a_1a_4 - a_2a_3 \geq 0.$$

By means of successive exterior differentiations of (1.5), we get the existence of functions $b_1, \dots, b_6, c_1, \dots, c_8$ such that

$$(1.7) \quad (da_1 - a_2\omega_1^2 - a_3\tau_1^2) \wedge \omega^1 + (da_2 + a_1\omega_1^2 - a_4\tau_1^2) \wedge \omega^2 = 0$$

$$(da_3 - a_4\omega_1^2 + a_1\tau_1^2) \wedge \omega^1 + (da_4 + a_3\omega_1^2 + a_2\tau_1^2) \wedge \omega^2 = 0;$$

$$(1.8) \quad da_1 - a_2\omega_1^2 - a_3\tau_1^2 = b_1\omega^1 + b_2\omega^2,$$

$$da_3 - a_4\omega_1^2 + a_1\tau_1^2 = b_4\omega^1 + b_5\omega^2,$$

$$da_2 + a_1\omega_1^2 - a_4\tau_1^2 = b_2\omega^1 + b_3\omega^2,$$

$$da_4 + a_3\omega_1^2 + a_2\tau_1^2 = b_5\omega^1 + b_6\omega^2;$$

$$(1.9) \quad (db_1 - 2b_2\omega_1^2 - b_4\tau_1^2) \wedge \omega^1 + \{db_2 + (b_1 - b_3)\omega_1^2 - b_5\tau_1^2\} \wedge \omega^2 =$$

$$= (a_2K + a_3\mu K^*) \omega^1 \wedge \omega^2,$$

$$\{db_2 + (b_1 - b_3)\omega_1^2 - b_5\tau_1^2\} \wedge \omega^1 + (db_3 + 2b_2\omega_1^2 - b_6\tau_1^2) \wedge \omega^2 =$$

$$= (-a_1K + a_4\mu K^*) \omega^1 \wedge \omega^2,$$

$$(db_4 - 2b_5\omega_1^2 + b_1\tau_1^2) \wedge \omega^1 + \{db_5 + (b_4 - b_6)\omega_1^2 + b_2\tau_1^2\} \wedge \omega^2 =$$

$$= (a_4K - a_1\mu K^*) \omega^1 \wedge \omega^2,$$

$$\{db_5 + (b_4 - b_6)\omega_1^2 + b_2\tau_1^2\} \wedge \omega^1 + (db_6 + 2b_5\omega_1^2 + b_3\tau_1^2) \wedge \omega^2 =$$

$$= (-a_3K - a_2\mu K^*) \omega^1 \wedge \omega^2;$$

$$(1.10) \quad db_1 - 2b_2\omega_1^2 - b_4\tau_1^2 = c_1\omega^1 + c_2\omega^2,$$

$$db_2 + (b_1 - b_3)\omega_1^2 - b_5\tau_1^2 = (c_2 + a_2K + a_3\mu K^*) \omega^1 +$$

$$+ (c_3 + a_1K - a_4\mu K^*) \omega^2,$$

$$db_3 + 2b_2\omega_1^2 - b_6\tau_1^2 = c_3\omega^1 + c_4\omega^2,$$

$$db_4 - 2b_5\omega_1^2 + b_1\tau_1^2 = c_5\omega^1 + c_6\omega^2,$$

$$db_5 + (b_4 - b_6)\omega_1^2 + b_2\tau_1^2 = (c_6 + a_4K - a_1\mu K^*) \omega^1 +$$

$$+ (c_7 + a_3K + a_2\mu K^*) \omega^2,$$

$$db_6 + 2b_5\omega_1^2 + b_3\tau_1^2 = c_7\omega^1 + c_8\omega^2.$$

Of course, we have

$$(1.11) \quad ds^2 = (\omega^1)^2 + (\omega^2)^2,$$

$$ds_*^2 = (\tau^1)^2 + (\tau^2)^2 =$$

$$= (a_1^2 + a_3^2)(\omega^1)^2 + 2(a_1a_2 + a_3a_4)\omega^1\omega^2 + (a_2^2 + a_4^2)(\omega^2)^2.$$

The fundamental invariants of f are

$$(1.12) \quad I_1 = a_1^2 + a_2^2 + a_3^2 + a_4^2, \quad I_2 = (a_1 - a_4)^2 + (a_2 + a_3)^2.$$

The mapping f is called *constant* if $I_1 = 0$; f is said to be *conformal* if $I_2 = 0$; the geometrical signification is obvious. To each point $m \in M$, we get the induced quadratic mapping

$$(1.13) \quad f_{**} : T_m(M) \rightarrow T_{f(m)}(N),$$

$$f_{**}(xv_1 + yv_2) = (b_1x^2 + 2b_2xy + b_3y^2)v_1^* + (b_4x^2 + 2b_5xy + b_6y^2)v_2^*;$$

see [2]. Further, we get the mapping

$$(1.14) \quad t : M \rightarrow T(N), \quad t(m) \in T_{f(m)}(N); \quad t = (b_1 + b_3)v_1^* + (b_2 + b_4)v_2^*;$$

t is the so-called *tension field*. The expressions

$$(1.15) \quad J_1 = (b_1 + b_3)^2 + (b_4 + b_6)^2, \quad J_2 = b_1^2 + 2b_2^2 + b_3^2 + b_4^2 + 2b_5^2 + b_6^2$$

are invariants of f as well; f is said to be *harmonic* if $J_1 = 0$, and it is *totally geodesic* if $J_2 = 0$.

2. Let us produce several integral formulas.

First of all, consider the 1-form

$$(2.1) \quad \varphi_1 = \{(a_1 - a_4)(b_2 + b_4) - (a_2 + a_3)(b_1 - b_5)\} \omega^1 +$$

$$+ \{(a_1 - a_4)(b_3 + b_5) - (a_2 + a_3)(b_2 - b_6)\} \omega^2;$$

φ_1 is invariant. Then

$$(2.2) \quad \int_{\partial M} \varphi_1 = \int_M \{2L_1 - I_2(K + \mu K^*)\} \omega^1 \wedge \omega^2,$$

$$\text{where } L_1 = \begin{vmatrix} b_1 - b_5 & b_2 - b_6 \\ b_2 + b_4 & b_3 + b_5 \end{vmatrix}.$$

For the invariant form

$$(2.3) \quad \varphi_2 = \{(a_1 + a_4)(b_2 - b_4) - (a_2 - a_3)(b_1 + b_5)\} \omega^1 +$$

$$+ \{(a_1 + a_4)(b_3 - b_5) - (a_2 - a_3)(b_2 + b_6)\} \omega^2,$$

we get

$$(2.4) \quad \int_{\partial M} \varphi_2 = \int_M \{2L_2 + I_3(\mu K^* - K)\} \omega^1 \wedge \omega^2,$$

$$\text{where } L_2 = \begin{vmatrix} b_1 + b_5 & b_2 + b_6 \\ b_2 - b_4 & b_3 - b_5 \end{vmatrix}$$

and

$$(2.5) \quad I_3 = (a_1 + a_4)^2 + (a_2 - a_3)^2 .$$

Further,

$$(2.6) \quad \frac{1}{2} dI_1 = (a_1 b_1 + a_2 b_2 + a_3 b_4 + a_4 b_5) \omega^1 + \\ + (a_1 b_2 + a_2 b_3 + a_3 b_5 + a_4 b_6) \omega^2 ,$$

i.e.,

$$(2.7) \quad \frac{1}{2} \int_{\partial M} * dI_1 = \int_M \{a_1(c_1 + c_3) + a_2(c_2 + c_4) + a_3(c_5 + c_7) + a_4(c_6 + c_8) + \\ + J_2 + I_1 K - 2\mu^2 K^*\} \omega^1 \wedge \omega^2 .$$

Analogously,

$$(2.8) \quad \frac{1}{2} dI_2 = \{(a_1 - a_4)(b_1 - b_5) + (a_2 + a_3)(b_2 + b_4)\} \omega^1 + \\ + \{(a_1 - a_4)(b_2 - b_6) + (a_2 + a_3)(b_3 + b_5)\} \omega^2$$

and

$$(2.9) \quad \frac{1}{2} \int_{\partial M} * dI_2 = \int_M \{(a_1 - a_4)(c_1 + c_3 - c_6 - c_8) + (a_2 + a_3)(c_2 + c_4 + c_5 + c_7) + \\ + (b_1 - b_5)^2 + (b_2 + b_4)^2 + (b_2 - b_6)^2 + (b_3 + b_5)^2 + I_2(K + \mu K^*)\} \omega^1 \wedge \omega^2 .$$

Next,

$$(2.10) \quad \frac{1}{2} dI_3 = \{(a_1 + a_4)(b_1 + b_5) + (a_2 - a_3)(b_2 - b_4)\} \omega^1 + \\ + \{(a_1 + a_4)(b_2 + b_6) + (a_2 - a_3)(b_3 - b_5)\} \omega^2$$

and

$$(2.11) \quad \frac{1}{2} \int_{\partial M} * dI_3 = \int_M \{(a_1 + a_4)(c_1 + c_3 + c_6 + c_8) + \\ + (a_2 - a_3)(c_2 + c_4 - c_5 - c_7) + (b_1 + b_5)^2 + (b_2 - b_4)^2 + \\ + (b_2 + b_6)^2 + (b_3 - b_5)^2 + I_3(K - \mu K^*)\} \omega^1 \wedge \omega^2 .$$

Finally, consider the invariant 1-form

$$(2.12) \quad \varphi_3 = \{(b_1 + b_3)(c_5 + c_7) - (b_4 + b_6)(c_1 + c_3)\} \omega^1 + \\ + \{(b_1 + b_3)(c_6 + c_8) - (b_4 + b_6)(c_2 + c_4)\} \omega^2 .$$

From (1.10), we get

$$(2.13) \quad \begin{aligned} d(b_1 + b_3) - (b_4 + b_6) \tau_1^2 &= (c_1 + c_3) \omega^1 + (c_2 + c_4) \omega^2, \\ d(b_4 + b_6) + (b_1 + b_3) \tau_1^2 &= (c_5 + c_7) \omega^1 + (c_6 + c_8) \omega^2. \end{aligned}$$

The exterior differentiation yields

$$(2.14) \quad \begin{aligned} &\{d(c_1 + c_3) - (c_2 + c_4) \omega_1^2 - (c_5 + c_7) \tau_1^2\} \wedge \omega^1 + \\ &+ \{d(c_2 + c_4) + (c_1 + c_3) \omega_1^2 - (c_6 + c_8) \tau_1^2\} \wedge \omega^2 = (b_4 + b_6) \mu K^* \omega^1 \wedge \omega^2, \\ &\{d(c_5 + c_7) - (c_6 + c_8) \omega_1^2 + (c_1 + c_3) \tau_1^2\} \wedge \omega^1 + \\ &+ \{d(c_6 + c_8) + (c_5 + c_7) \omega_1^2 + (c_2 + c_4) \tau_1^2\} \wedge \omega^2 = -(b_1 + b_3) \mu K^* \omega^1 \wedge \omega^2 \end{aligned}$$

and the existence of functions e_1, \dots, e_6 such that

$$(2.15) \quad \begin{aligned} d(c_1 + c_3) - (c_2 + c_4) \omega_1^2 - (c_5 + c_7) \tau_1^2 &= e_1 \omega^1 + (e_2 - b_4 \mu K^*) \omega^2, \\ d(c_2 + c_4) + (c_1 + c_3) \omega_1^2 - (c_6 + c_8) \tau_1^2 &= (e_2 + b_6 \mu K^*) \omega^1 + e_3 \omega^2, \\ d(c_5 + c_7) - (c_6 + c_8) \omega_1^2 + (c_1 + c_3) \tau_1^2 &= e_4 \omega^1 + (e_5 + b_1 \mu K^*) \omega^2, \\ d(c_6 + c_8) + (c_5 + c_7) \omega_1^2 + (c_2 + c_4) \tau_1^2 &= (e_5 - b_3 \mu K^*) \omega^1 + e_6 \omega^2. \end{aligned}$$

By means of (2.13) and (2.15), we get the integral formula

$$(2.16) \quad \int_{\partial M} \varphi_3 = \int_M (2L_3 - J_1 \mu K^*) \omega^1 \wedge \omega^2, \quad \text{where } L_3 = \begin{vmatrix} c_1 + c_3 & c_2 + c_4 \\ c_5 + c_7 & c_6 + c_8 \end{vmatrix}.$$

3. Let us explain the geometrical interpretation of the invariants L_i . Introduce the invariant operator

$$(3.1) \quad * : T_m(M) \rightarrow T_m(M), \quad m \in M; \quad *(xv_1 + yv_2) = -yv_1 + xv_2;$$

satisfying $\omega(*v) + *\omega(v) = 0$ for $v \in T_m(M)$, $\omega \in T_m^*(M)$. To f_{**} (1.13), consider the associated bilinear mapping

$$(3.2) \quad \begin{aligned} \mathcal{L} : T_m(M) \times T_m(M) &\rightarrow T_{f(m)}(N), \\ \mathcal{L}(x^1 v_1 + x^2 v_2, y^1 v_1 + y^2 v_2) &= (b_1 x^1 y^1 + b_2 x^1 y^2 + b_2 x^2 y^1 + b_3 x^2 y^2) v_1^* + \\ &+ (b_4 x^1 y^1 + b_5 x^1 y^2 + b_5 x^2 y^1 + b_6 x^2 y^2) v_2^*. \end{aligned}$$

Finally, consider the operator

$$(3.3) \quad * : T_n(N) \rightarrow T_n(N), \quad n \in N; \quad *(\xi v_1^* + \eta v_2^*) = -\eta v_1^* + \xi v_2^*.$$

Lemma 1. Let $v \in T_m(M)$ be an arbitrary unit vector. Then

$$(3.4) \quad \begin{aligned} L_1 &= \langle \mathcal{L}(v, v) + *\mathcal{L}(v, *v), \mathcal{L}(*v, *v) - *\mathcal{L}(v, *v) \rangle, \\ L_2 &= \langle \mathcal{L}(v, v) - *\mathcal{L}(v, *v), \mathcal{L}(*v, *v) + *\mathcal{L}(v, *v) \rangle. \end{aligned}$$

Proof. Because of the invariance of L_1, L_2 and \mathcal{L} , we may choose the frames such that $v = v_1$, i.e., $*v = v_2$ at $m \in M$. Then

$$\begin{aligned} \mathcal{L}(v_1, v_1) &= b_1 v_1^* + b_4 v_2^*, \\ \mathcal{L}(v_1, v_2) &= b_2 v_1^* + b_5 v_2^*, \quad \mathcal{L}(v_2, v_2) = b_3 v_1^* + b_6 v_2^*, \end{aligned}$$

and our Lemma follows. QED.

Lemma 2. Let t (1.14) be the tension field, $v \in T_m(M)$ an arbitrary unit vector and

$$(3.5) \quad V := f_* v, \quad W := f_*(*v).$$

Then

$$(3.6) \quad L_3 = \langle *\nabla_V^* t, \nabla_W^* t \rangle.$$

Proof. We have

$$\begin{aligned} \nabla^* t &= (b_4 + b_6)(\tau_1^2 - \Omega_1^2)v_1^* + (b_1 + b_3)(\Omega_1^2 - \tau_1^2)v_2^* + \\ &+ \{(c_1 + c_3)\omega^1 + (c_2 + c_4)\omega^2\}v_1^* + \{(c_5 + c_7)\omega^1 + (c_6 + c_8)\omega^2\}v_2^*. \end{aligned}$$

Notice that, for each form $\Omega \in T^*(N)$ and each vector $v \in T(M)$, we have $f^* \Omega(v) = \Omega(f_* v)$. Set $v = v_1$ at $m \in M$. Then

$$\nabla_V^* t = (c_1 + c_3)v_1^* + (c_5 + c_7)v_2^*, \quad \nabla_W^* t = (c_2 + c_4)v_1^* + (c_6 + c_8)v_2^*,$$

and the Lemma follows. QED.

Lemma 3. Let t (1.14) be the tension field and V, W be defined by (3.5). Let $\varepsilon = \pm 1$. Then

$$(3.7) \quad \nabla_V^* t + \varepsilon * \nabla_W^* t = 0$$

for each $v \in T(M)$ if and only if

$$(3.8) \quad c_1 + c_3 - \varepsilon(c_6 + c_8) = c_2 + c_4 + \varepsilon(c_5 + c_7) = 0.$$

Proof. Let $v = xv_1 + yv_2$. Then

$$\begin{aligned} \nabla_V^* t &= \{(c_1 + c_3)x + (c_2 + c_4)y\}v_1^* + \{(c_5 + c_7)x + (c_6 + c_8)y\}v_2^*, \\ \nabla_W^* t &= \{(c_2 + c_4)x - (c_1 + c_3)y\}v_1^* + \{(c_6 + c_8)x - (c_5 + c_7)y\}v_2^* \end{aligned}$$

and

$$\begin{aligned} \nabla_{\bar{v}}^* t + \varepsilon * \nabla_{\bar{w}}^* t = & \{c_1 + c_3 - \varepsilon(c_6 + c_8)\} (xv_1^* - \varepsilon yv_2^*) + \\ & + \{c_2 + c_4 + \varepsilon(c_5 + c_7)\} (yv_1^* + \varepsilon xv_2^*). \end{aligned}$$

The Lemma follows easily. QED.

4. Our main task is to obtain several typical geometric consequences of our integral formulas (2.2, 4, 7, 9, 11, 16). In all theorems, M and N are Riemannian manifolds, $\dim M = \dim N = 2$, $f: M \rightarrow N$ is an orientation preserving mapping, ∂M the boundary of M . All other notions have been explained above.

First of all, let us state the following

Lemma 4. *The condition*

$$(4.1) \quad f \text{ is harmonic}$$

and/or the condition

$$(4.2) \quad \text{for each } m \in M \text{ there is } \dim f_{**}(T_m(M)) \leq 1 \text{ and there exists a vector } 0 \neq v \in T_m(M) \text{ such that } f_{**}(v) = 0$$

implies

$$(4.3) \quad L_1 \leq 0, \quad L_2 \leq 0.$$

Proof. The condition (4.1) is equivalent to $b_1 + b_3 = b_4 + b_6 = 0$. Hence

$$L_1 = -(b_1 - b_5)^2 - (b_2 + b_4)^2 \leq 0, \quad L_2 = -(b_1 + b_2)^2 - (b_2 - b_4)^2 \leq 0.$$

From (4.2), we get — see (1.13) — the existence of functions $\varrho, \sigma, B_1, B_2, B_3$ such that

$$b_1 = \varrho B_1, \quad b_2 = \varrho B_2, \quad b_3 = \varrho B_3, \quad b_4 = \sigma B_1, \quad b_5 = \sigma B_2, \quad b_6 = \sigma B_3$$

and

$$f_{**}(xv_1 + yv_2) = (B_1x^2 + 2B_2xy + B_3y^2)(\varrho v_1^* + \sigma v_2^*).$$

Further,

$$L_1 = L_2 = (\varrho^2 + \sigma^2)(B_1B_3 - B_2^2),$$

and our Lemma follows. QED.

Theorem 1. *Suppose: (i) $L_1 \leq 0$ on M , (ii) $K > 0$ on M , (iii) $K^* \geq 0$ on $f(M) \subset N$, (iv) $I_2 = 0$ on ∂M . Then f is conformal.*

Proof is a direct consequence of (2.2). QED.

Theorem 2. *Suppose: (i) M is compact, (ii) $L_2 \leq 0$ on M , (iii) $K > 0$ on M ,*

(iv) $K^* \leq 0$ on $f(M) \subset N$. Then f is a constant mapping. We may suppose (i') $J_2 = 0$ on ∂M instead of (i).

Proof. From (2.4), $a_1 + a_4 = a_2 - a_3 = 0$. Hence $\mu = -a_1^2 - a_2^2$; from $\mu \geq 0$, we get $a_1 = a_2 = 0$. QED.

Theorem 3. Suppose: (i) f is harmonic, (ii) $K \geq 0$ on M , (iii) $K^* \leq 0$ on $f(M) \subset N$, (iv) $J_2 = 0$ on ∂M . Then f is totally geodesic. Replacing (ii) by (ii') $K > 0$ on M , f has to be a constant mapping.

Proof is a consequence of (2.7). QED.

Theorem 4. Suppose: (i) for each $v \in T(M)$, we have $\nabla_v^* t + * \nabla_w^* t = 0$, t being the tension field and $V := f_* v$, $W := f_*(v)$, (ii) $K > 0$ on M , (iii) $K^* \geq 0$ on $f(M) \subset N$, (iv) $I_2 = 0$ on ∂M . Then f is conformal.

Proof follows from Lemma 3 and (2.9). QED.

Theorem 5. Suppose: (i) for each $v \in T(M)$, we have $\nabla_v^* t + * \nabla_w^* t = 0$, t being the tension field and $V := f_* v$, $W := f_*(v)$, (ii) $K \geq 0$ on M , (iii) $K^* \geq 0$ on $f(M) \subset N$, (iv) $I_2 = 0$ on ∂M . Then f is harmonic and we have

$$(4.4) \quad I_2(K + \mu K^*) = 0$$

at each point $m \in M$.

Proof. From Lemma 3 and (2.9),

$$b_1 - b_5 = b_2 + b_4 = b_2 - b_6 = b_3 + b_5 = 0$$

and f is to be harmonic. From (1.10),

$$\begin{aligned} c_1 - c_6 - a_4 K + a_1 \mu K^* &= 0, & c_2 - c_7 + a_2 K + a_3 \mu K^* &= 0, \\ c_2 - c_7 - a_3 K - a_2 \mu K^* &= 0, & c_3 - c_8 + a_1 K - a_4 \mu K^* &= 0, \\ c_2 + c_5 + a_2 K + a_3 \mu K^* &= 0, & c_3 + c_6 + a_4 K - a_1 \mu K^* &= 0, \\ c_3 + c_6 + a_1 K - a_4 \mu K^* &= 0, & c_4 + c_7 + a_3 K + a_2 \mu K^* &= 0. \end{aligned}$$

By the elimination of c_1, \dots, c_8 from these equations and from (3.8) for $\varepsilon = 1$, we get

$$(a_1 - a_4)(K + \mu K^*) = (a_2 + a_3)(K + \mu K^*) = 0,$$

i.e., (4.4). QED.

Theorem 6. Suppose: (i) M is compact, (ii) for each $v \in T(M)$, we have $\nabla_v^* t = * \nabla_w^* t$, t being the tension field and $V := f_*(v)$, $W := f_*(v)$, (iii) $K > 0$ on M ,

(iv) $K^* \leq 0$ on $f(M) \subset N$. Then f is a constant mapping. Instead of (i), it is sufficient to suppose (i') $J_2 = 0$ on ∂M .

Proof. From Lemma 3 and (2.11), we get $I_3 = 0$, i.e., $\mu = -a_1^2 - a_2^2$. From $\mu \geq 0$, we get $a_1 = a_2 = 0$. QED.

Theorem 7. Suppose: (i) M is compact, (ii) for each $v \in T(M)$, $\nabla_v^* t = * \nabla_W^* t$, t being the tension field and $V := f_*(v)$, $W := f_*(v)$, (iii) $K \geq 0$ on M , (iv) $K^* \leq 0$ on $f(M) \subset N$. Then f is harmonic and we have

$$(4.5) \quad I_3(K - \mu K^*) = 0$$

at each point $m \in M$. Instead of (i), it is sufficient to suppose (i') $J_2 = 0$ on ∂M .

Proof. From Lemma 3 and (2.11),

$$b_1 + b_5 = b_2 - b_4 = b_2 + b_6 = b_3 - b_5 = 0,$$

and f is harmonic. From (1.10),

$$\begin{aligned} c_1 + c_6 + a_4 K - a_1 \mu K^* &= 0, & c_2 + c_7 + a_2 K + a_3 \mu K^* &= 0, \\ c_2 + c_7 + a_3 K + a_2 \mu K^* &= 0, & c_3 + c_8 + a_1 K - a_4 \mu K^* &= 0, \\ c_2 - c_5 + a_2 K + a_3 \mu K^* &= 0, & c_3 - c_5 - a_4 K + a_1 \mu K^* &= 0, \\ c_3 - c_6 + a_1 K - a_4 \mu K^* &= 0, & c_4 - c_7 - a_3 K - a_2 \mu K^* &= 0. \end{aligned}$$

The elimination of c_1, \dots, c_8 from these equations and from (3.8) for $\varepsilon = -1$ implies

$$(a_1 + a_4)(K - \mu K^*) = 0, \quad (a_2 - a_3)(K - \mu K^*) = 0,$$

i.e., (4.5). QED.

Theorem 8. Suppose: (i) $L_3 \geq 0$ on M , (ii) $K^* < 0$ on $f(M) \subset N$, (iii) $f_*(T_m(M)) = T_{f(m)}(N)$ for each $m \in M$, (iv) $J_1 = 0$ on ∂M . Then f is a harmonic mapping.

Proof. From (2.16), $J_1 = 0$ on M . QED.

Theorem 9. Suppose: (i) M is compact, (ii) $J_1 = \text{const.} \neq 0$ on M , (iii) $K^* > 0$ or $K^* < 0$ on $f(M) \subset N$. Then $\dim f_*(T_m(M)) \leq 1$ for each $m \in M$.

Proof. From $J_1 = \text{const.}$ and (1.10), we get

$$\begin{aligned} (b_1 + b_3)(c_1 + c_3) + (b_4 + b_6)(c_5 + c_7) &= 0, \\ (b_1 + b_3)(c_2 + c_4) + (b_4 + b_6)(c_6 + c_8) &= 0 \end{aligned}$$